VIBRANT MEscope Application Note 28

Mathematics of a Mass-Spring-Damper System

The steps in this Application Note can be carried out using any MEscope package that includes the **VES-3600 Advanced Signal Processing** and **VES-4000 Modal Analysis** options. Without these options, you can still carry out the steps in this App Note using the **AppNote28** project file. These steps might also require a *more recent release date* of MEscope.

APP NOTE 28 PROJECT FILE

• To retrieve the Project file for this App Note, click here to download AppNote28.zip

This Project contains numbered Hotkeys & Scripts of commands for carrying out the steps of this App Note.

• Hold down the Ctrl key and click on a Hotkey to open its Script window

INTRODUCTION

In this note, MEscope is used to explore the properties of the mass-spring-damper system shown in the figure below. Its equation of motion will be solved for its mode of vibration. Then, an **FRF** will be synthesized using its mode shape, and its stiffness and mass lines will be examined.

Then, the **FRF** will be curve fit to extract its modal parameters. Finally, we will look at how the modal parameters are contained in the Impulse Response Function, the Inverse FFT of the **FRF**.



Mass-Spring-Damper.

The purpose of this Application Note is to review the details of modal analysis, to provide a better understanding of the modal properties of all structures. The modal properties of real-world structures are analyzed using a multi-degree-of-freedom (MDOF) dynamic model, whereas the model used here is a single degree-of-freedom (**SDOF**) model. Nevertheless, the dynamics of MDOF structures are better understood by analyzing the dynamics of this **SDOF** structure.

Modes are defined for structures, the dynamics of which can be represented by *linear ordinary differential equations* like the one in the **Background Math** section below. The dynamic behavior of the mass-spring-damper structure in the figure above is represented by a single (scalar) equation. An MDOF structure is represented by multiple equations, which are written in matrix form.

Because of the *superposition property* of linear systems, the dynamics of an MDOF structure can be written as a *summation of contributions due to each of its mode shapes*. Each mode shape can be thought of as representing the dynamics of a single mass-spring-damper system.



Mass-Spring-Damper 3D Model.

BACKGROUND MATH

The time domain equation of motion for the mass-spring-damper is represented by *Newton's Second Law*, written as the following *force balance* on a structure,

$$\mathbf{M}\,\ddot{\mathbf{x}}(\mathbf{t}) + \mathbf{C}\,\dot{\mathbf{x}}(\mathbf{t}) + \mathbf{K}\,\mathbf{x}(\mathbf{t}) = \mathbf{f}(\mathbf{t})$$

- M → Mass
- C → Damping
- $\mathbf{K} \rightarrow \mathbf{Stiffness}$
- $\ddot{\mathbf{x}}(\mathbf{t}) \rightarrow \text{acceleration}$
- $\dot{\mathbf{x}}(\mathbf{t}) \rightarrow \text{velecity}$
- $\mathbf{x}(\mathbf{t}) \rightarrow \text{displacement}$
- $f(t) \rightarrow$ excitation force

LAPLACE TRANSFORMS

By taking Laplace transforms of the terms in the differential equation above and setting initial conditions to zero, an equivalent frequency domain equation of motion results,

> [M s² + C s + K] X(s) = F(s)X(s) → Laplace transform of the displacement F(s) → Laplace transform of the force $s = \sigma + j\omega$ → complex Laplace variable

TRANSFER FUNCTION

The above equation can be rewritten by simply dividing both sides by the coefficients of the left-hand side.

$$X(s) = \left(\frac{1}{Ms^2 + C s + K}\right) F(s)$$

The new coefficient on the right-hand side of the above equation is called the Transfer Function,

$$\mathbf{H}(\mathbf{s}) = \frac{\mathbf{X}(\mathbf{s})}{\mathbf{F}(\mathbf{s})} = \left(\frac{1}{\mathbf{M}\mathbf{s}^2 + \mathbf{C}\mathbf{s} + \mathbf{K}}\right)$$

The Transfer Function is *complex valued* and has *real & imaginary* parts or equivalently *magnitude & phase*. The two parts of the Transfer Function can be plotted on the complex Laplace plane (or **s**-plane), as shown below.



Transfer Function Plotted Over Half of the S-Plane.

POLES OF THE TRANSFER FUNCTION

Notice that the magnitude of the Transfer Function *has two peaks*. These are points where the value of the Transfer Function *goes to infinity*. The *real & imaginary parts* also have the *same two peaks*.

The Transfer Function goes to infinity for values on the **s**-Plane where its *denominator is zero*. It is also clear that *as s goes to infinity*, the Transfer Function *approaches zero*.

CHARACTERISTIC POLYNOMIAL

The denominator of the Transfer Function is a *second order polynomial* in the **s** variable and is called the *characteristic polynomial*. Since it is a second order polynomial, it has two roots (values of **s** for which *it is zero*).

These two roots of the denominator are called the *poles* of the Transfer Function.

The poles are *complex conjugates* of one another. The poles are the locations on the **s**-plane where the Transfer Function has a *value of infinity*. The poles are also called *eigenvalues*.

 $p_0=-\sigma_0+j\omega_0$, $p_0^*=-\sigma_0-j\omega_0$

S-PLANE NOMENCLATURE

The real axis in the S-Plane is called the damping axis, and the imaginary axis is called the frequency axis.

The locations of the poles in the S-Plane have also been given some other commonly used names, as shown below.



S-Plane Nomenclature.

MODAL PARAMETERS

The coordinates of the poles in the **s**-Plane are also modal parameters. The Transfer Function can be written in terms of its pole locations, or modal parameters,

$$H(s) = \left(\frac{1/M}{s^2 + 2\sigma_0 s + \Omega_0^2}\right)$$
$$\sigma_0 = \frac{C}{2 M}, \Omega_0^2 = \frac{K}{M}$$
$$\Omega_0^2 = \sigma_0^2 + \omega_0^2$$
$$\sigma_0 \Rightarrow \text{ modal damping}$$
$$\Omega_0 \Rightarrow \text{ un-damped modal frequency}$$
$$\omega_0 \Rightarrow \text{ damped modal frequency}$$
The **percent of critical damping** (\zeta_0) is written as,

$$\zeta_0 = \frac{\sigma_0}{\Omega_0} = \frac{C}{2\sqrt{MK}}$$

FREQUENCY RESPONSE FUNCTION (FRF)

In the figure below, *the Transfer Function has only been plotted for half of the s-Plane*. It has only been plotted for *negative values of* σ (the real part of s). This was done so that the values of the Transfer Function along the j ω -axis (the imaginary part of s) are clearly seen.

The Frequency Response Function (**FRF**) is the values of the Transfer Function *along the* j₀-axis.

The **FRF** values along the **j** $\boldsymbol{\omega}$ -axis are shown with black lines below. Since the **FRF** is only defined along the **j** $\boldsymbol{\omega}$ -axis, **s** can be replaced by its imaginary part **j** $\boldsymbol{\omega}$,

$$\begin{aligned} \mathbf{FRF} &= \mathbf{H}(\mathbf{j}\omega) = \mathbf{H}(\mathbf{s}) \Big|_{\mathbf{s} = \mathbf{j}\omega} \\ &= \frac{\mathbf{X}(\mathbf{s})}{\mathbf{F}(\mathbf{s})} \Big|_{\mathbf{s} = \mathbf{j}\omega} = \frac{\mathbf{X}(\mathbf{j}\omega)}{\mathbf{F}(\mathbf{j}\omega)} \end{aligned}$$



FRF Plotted on the jo-axis.

FRF IN PARTIAL FRACTION FORM

The **FRF** for an **SDOF** (mass-spring-damper) can now be written in terms of modal parameters by also replacing the **s**-variable with $j\omega$ in the equation for the Transfer Function in terms of modal parameters,

$$H(j\omega) = \frac{(1/M)}{(j\omega)^2 + 2\sigma_0 j\omega + {\Omega_0}^2}$$

Furthermore, using the poles of the Transfer Function, the FRF can be written in a partial fraction expansion form,

$$H(j\omega) = \frac{1}{2j} \left[\frac{R_0}{j\omega - p_0} - \frac{R_0}{j\omega - p_0^*} \right]$$

$$R_0 = 1/\omega_0 M$$

$$R_0 \Rightarrow \text{modal residue}$$

$$p_0 = -\sigma_0 + j\omega_0, p_0^* = -\sigma_0 - j\omega_0$$

The modal residue is the amplitude (or strength) of the numerator of each resonance term in the above equation.

The FRF of an SDOF is the summation of two resonance curves, each one with a peak near a pole location.

The Partial Fraction Form of the FRF for an **SDOF** is fully represented by two poles and two residues.

Since the residues are equal and the poles are complex conjugates of one another, the complete dynamics of an **SDOF** system is *fully represented by a modal frequency* $(j\omega_0)$, *modal damping* (σ_0) , and *modal residue* (\mathbf{R}_0) .

MODE SHAPES

One final step is to represent the FRF in terms of mode shapes instead of residues,

$$H(j\omega) = \frac{1}{2j} \left[\frac{\{u_0\}^2}{j\omega - p_0} - \frac{\{u_0\}^2}{j\omega - p_0^*} \right]$$
$$\{u_0\} = \left\{ \frac{1}{\sqrt{A\omega_0 M}} \right\}, A \rightarrow \text{scaling constant}$$

The mode shape $\{\mathbf{u}_0\}$ is a vector. Its first component is the *square root of the residue*, and its second component is zero. The second component *corresponds to the ground*, where there is no motion.

The mode shape also contains a scaling constant (A).

A mode shape *does not have unique values*. Only its *shape is unique*, one component relative to another.

Because it does not have unique values, a mode shape is also *called an eigenvector*.

The dynamics of an **SDOF** is fully *represented by two poles (eigenvalues)* and *two mode shapes (eigenvectors)*.

Because of symmetry, the dynamics of an SDOF *is fully represented* by a modal frequency $(j\omega_0)$, modal damping (σ_0) , and a mode shape $\{u_0\}$

IMPULSE RESPONSE FUNCTION (IRF)

An Impulse Response Function is the Inverse FFT of an FRF.

When it is written in terms of modal parameters, an **IRF** provides *the best source of meaning* for modal parameters.

In terms of modal parameters, the IRF is written,

$$h(t) = FFT^{-1} \left(\frac{1}{2j} \left[\frac{R_0}{j\omega - p_0} - \frac{R_0}{j\omega - p_0^*} \right] \right)$$
$$h(t) = \frac{T}{2j} \left[R_0 e^{p_0 t} - R_0 e^{p_0^* t} \right]$$
$$h(t) = T |R_0| e^{-\sigma_0 t} (\sin(\omega_0 t + \alpha_0))$$
$$\alpha_0 \rightarrow \text{angle of } R_0$$

The equation above shows the role that each modal parameter plays in the IRF.

 $\omega_0 \rightarrow$ multiplies the time variable (t) in the sinusoidal function $(\sin(\omega_0 t + \alpha_0))$

 $\omega_0 \rightarrow$ defines the frequency of oscillation, hence ω_0 is called *damped modal frequency*

 $\sigma_0 \rightarrow$ the coefficient in the exponential term $(e^{-\sigma_0 t})$ that defines the *envelope of decay* of the IRF

 $\sigma_0 \Rightarrow$ called the *damping decay constant*, *modal damping coefficient*, *half power point damping*, 3 *dB point damping*

In a real-world structure, vibration decay is caused by a *combination of damping mechanisms* within or about the structure. On earth, the *most common damping mechanism* is the *surrounding air*.

FEA MODEL

In MEscope, a graphical 3D model is needed to display experimentally derived ODS's and mode shapes in animation. The mass-spring damper model above was created in **App Note 07A** using *five substructures*. Building a model using substructures makes setting up the model for shape animation and FEA modeling much easier.

In this App Note, an FEA mass, spring, & damper *have already been added* to the 3D model.

The resulting FEA model will be solved for its single mode of vibration with and without the damper included.

The FEA models of most real-world structures *do not include damping* because *it is too difficult* to include an accurate damping model of the structure into a set of MDOF (multiple degree-of-freedom) differential equations. When damping is not included, the FEA model can still be solved for its mode shapes.

Mode shapes of an **undamped structure** are called **normal modes**.

FEA MASS, STIFFNESS & DAMPING

The FEA properties of the Mass, Spring & Damper have already been added to the FEA Properties dialog. To display the FEA Properties,

• Open an STR window and execute FEA | FEA Properties

The following properties are displayed on the Springs, Dampers, and Masses tabs.

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MODAL FREQUENCY & DAMPING

Millimeter (mm) units are internally converted to Meters (m) to make them consistent with Newtons (N) and kilograms (kg).

The un-damped natural frequency (f_n) of the Mass-Spring-Damper is calculated with the following formula

$$f_n = \frac{1}{2\pi} \sqrt{\frac{K}{M}} = \frac{1}{2\pi} \sqrt{\frac{2000}{0.5}} = 10.07 \text{ Hz}$$

The **percent of critical damping** (ζ) is calculated with the following formula,

$$\zeta = \frac{C}{\sqrt{4KM}} = \frac{0.005}{\sqrt{4 \times 2000 \times 0.5}} = 7.906 \%$$

The damped natural frequency (f_d) is calculated with the following formula,

$$f_d = f_n \sqrt{1-\zeta^2} = 10.03 \text{ Hz}$$



STEP 1 – DAMPED AND UNDAMPED MODE SHAPE

• *Press* Hotkey 1 Damp vs. Undamped Mode Shape

In this step, the FEA mode shape of the **damped SDOF model** is compared with the mode shape of the **undamped SDOF model**.

The Mass, Spring & Damper have already been attached to the 3D model

The Mass is attached to a Point in the center of the mass cube

The Spring & Damper have been attached between the Mass Point and a Fixed base Point. All the Points on the Base plate are Fixed

After the undamped and damped FEA mode shapes have been calculated, they are displayed side-by-side in animation, as shown below.



Damped & Undamped Mode Shapes Side-by-Side.

MAC \rightarrow 1.0 between the two mode shapes, meaning that *they are identical*. That is also apparent by observing the **M#s** in each Shape Table *on the right*. The FEA solver also calculated rotational DOFs but those are not displayed in animation of the **SDOF** models.

The modal frequencies, damped \rightarrow 10.03 Hz, undamped \rightarrow 10.07 Hz, are identical to the values calculated with the analytical formulas.

UMM MODE SHAPES

In the two Shape Tables listed above, the M#s of the FEA mode shapes are designated as UMM Mode Shapes.

In MEscope mode shapes scaled to Unit Modal Masses are called **UMM** mode shapes.

A common way to scale mode shapes is so that they yield unit modal masses. If the mass is pre- and post-multiplied by a **UMM** mode shape, the **resulting modal mass** \rightarrow 1. For an MDOF system, if the mass matrix is pre- and post-multiplying by a set of **UMM** mode shapes the resulting modal mass matrix \rightarrow 1 *along the diagonal*.

When FEA mode shapes are calculated in MEscope, the resulting mode shapes are **UMM** mode shapes.

THE DRIVING POINT FRF

An **FRF**, like a Transfer Function, defines the dynamic characteristics between two DOFs of a structure. MDOF systems have many DOF pairs for which FRFs can be calculated. The **SDOF** system in this App Note has only one meaningful driving point at **DOF 39Z**. Force can be applied at the mass point (**Point 39**) in the vertical (**Z**) direction. Using the damped **UMM** mode shape an **FRF** can be synthesized between the DOF (**39Z**) and itself.

Any FRF between a DOF and itself is called a driving point FRF.

Since an **FRF** is a ratio of response (**Output**) divided by excitation (**Input**). In this case, the units of the driving point **FRF** are (**mm/N**), and its DOFs are denoted **39Z:39Z**.

DRIVING POINT RESIDUE

UMM mode shapes are *scaled differently* than **Residue** mode shapes. For an **SDOF** model, the driving point Residue is equated to *the driving point component* of the **UMM mode shape** by the formula,

$$\mathbf{R}_0 = \frac{\{\mathbf{U}_0\}^2}{\mathbf{A}\boldsymbol{\varpi}_0\mathbf{M}}$$

A Residue mode shape has unique values and engineering units, in this case (mm/(N-sec))

A UMM mode shape has no unique values. Its shape is unique, but its values are not

Since the scaling constant (A) of a mode shape is arbitrary, it can always be equal to 1

With A=1, the driving point component of the UMM mode shapes in the previous figure can be scaled to driving point Residues

$$\mathbf{R}_0 = \frac{\{\mathbf{U}_0\}^2}{\mathbf{A}\omega_0 \mathbf{M}} = \frac{\{44.72\}^2}{((2\,\pi)\mathbf{10})(0.5)} \equiv 31.72$$

To verify this result,

• Execute Tools | Scaling | UMM to Residue Shapes and select DOF 39Z, the driving point DOF at the Mass

STEP 2 - SYNTHESIZING THE DRIVING POINT FRF

Press Hotkey 2 Driving Point FRF

The driving Point **FRF** will be displayed as shown below.



Driving Point FRF 39Z:39Z in (mm/N) Units.

STIFFNESS LINE OF THE FRF

As the frequency approaches **0 Hz**, the equation for the **FRF** shows that a (displacement/force) **FRF** approaches a constant, equal to the *inverse of the stiffness*.

The *inverse of the* stiffness is called the flexibility.

By letting **S=0** in the Transfer Function equation, the flexibility is,

$$H(0) \approx \frac{1}{K} = \frac{1}{2} = 0.5 \text{ (mm/N)}$$

• The flexibility is verified by the cursor value on the Log Magnitude in the figure above.

STEP 3 - MASS LINE OF THE FRF

• Press Hotkey 3 Mass Line of the FRF

As frequency becomes large, the **FRF** becomes dominated by the $j\omega^2$ term in its denominator, or characteristic polynomial. Therefore, at high frequencies the magnitude of the **FRF** can be approximated by,

$$|H(\omega)|\approx \left(\frac{1}{M\omega^2}\right)$$

When Hotkey 3 is *pressed*, the Tools | Differentiate command *is executed twice* to multiply the (displacement/force) FRF by ω^2 . Hence, its magnitude for high frequencies becomes a constant,

$$\omega^2 \ |H(\omega)| \approx \left(\frac{1}{M}\right)$$

The cursor value on the Log Magnitude below is *approximately* 2 (m/sec) $^2/N \rightarrow 0.5$ Kg, the mass of the Mass-Spring-Damper structure.

One Newton \rightarrow (one kilogram) x (one meter per second squared).



Line Cursor Value Showing the Value of the Mass Line.

STEP 4 - QUICK FIT OF THE FRF

• Press Hotkey 4 Quick fit of the FRF

A Driving Point **FRF** of the **SDOF** is a complete representation of the dynamics of the mass-spring-damper system at its driving point.

Curve fitting an **FRF** will recover the modal parameters of the **SDOF**, which are also *a complete representation of the driving point dynamics*.

The modal parameter estimates are displayed in the Modal Parameters spreadsheet shown below on the lower right.



Curve Fitting the Synthesized FRF.

The modal frequency, damping & residue values are all the correct values of the damped SDOF system.

IMPULSE RESPONSE FUNCTION (IRF)

The **IRF** is the product of two functions, an *exponential decay function* $(T | R_0 | e^{-\sigma_0 t})$ and a *sinusoidal function* $(sin(\omega_0 t + \alpha))$. The exponential decay $(e^{-\sigma_0 t})$ and sinusoidal function *are dimensionless*.

Only **T** and $\mathbf{R_0}$ have units

IRF units \rightarrow (Seconds) x (Residue units)

IRF units **→ FRF** units

ENVELOPE OF DECAY

The *maximum possible magnitude* of the IRF is the value of the exponential function for t=0.

max. mag. = T $|R_0|e^{-\sigma_0 t}|_{t=0} = T |R_0|$

This is the product of the time length (**T**) of the **IRF**, and the magnitude of the residue $|\mathbf{R}_0|$.

To check the envelope of the **IRF**, a new **FRF** will be synthesized using the modal parameters, but this time the **frequency increment** will be chosen so that T=1. One of the equations governing the FFT is,

$$\Delta f = \frac{1}{T} = 1$$
 Hz

Therefore, it an **FRF** is synthesized with $\Delta f \rightarrow 1$, and the *Inverse FFT* is applied to it, an **IRF** with **T=1** will result.

STEP 5 - LOG DECREMENT OF THE IRF

• Press Hotkey 5 Log Decrement of the IRF

The **IRF** shown below is displayed with the Line cursor at the maximum peak.

Peak value of the **IRF** \rightarrow 28.07 mm/N

Since T=1, the maximum envelope value (for t = 0) should be *approximately equal* to the *magnitude of the residue*,

R₀ → 31.72 mm/N-sec

Another Data Block with the value of the modal damping (in Hz) is displayed on the right.

This is the Log Decrement, calculated as the slope of the peaks in the the Log Magnitude of the IRF.

Log Decrement → 0.7887 → closely matches 0.7958 Hz



Log Decrement of the IRF.

STEP 6 - REVIEW STEPS

To review the steps of this App Note,

• Press Hotkey 6 Review Steps